

# Alternative formulations for statistical theories: 2.

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Carrying on from my previous post, I thought it would be interesting to look at the effect of the different formulations on statistical closure theories. In order to keep matters as simple as possible, I am restricting my attention to single-time theories and their forms for the transfer spectrum  $T(k,t)$  as it occurs in the Lin equation (see page 56 in [1]). For instance, the form for this due to Edwards [2] may be written in terms of the spectral energy density  $C(k,t)$  (or spectral covariance) as:

$$T(k,t) = 4\pi \int d^3j L(k,j,|\mathbf{k}-\mathbf{j}|) D(k,j,|\mathbf{k}-\mathbf{j}|) C(|\mathbf{k}-\mathbf{j}|,t) [C(j,t) - C(k,t)],$$

where

$$D(k,j,|\mathbf{k}-\mathbf{j}|) = \frac{1}{\omega(k,t) + \omega(j,t) + \omega(|\mathbf{k}-\mathbf{j}|,t)},$$

and  $\omega(k,t)$  is the inverse modal response time. The geometric factor  $L(\mathbf{k},\mathbf{j})$  is given by:

$$L(\mathbf{k},\mathbf{j}) = \frac{[\mu(k^2+j^2) - kj(1+2\mu^2)]}{(1-\mu^2)kj(k^2+j^2-2kj\mu)},$$

and can be seen by inspection to have the symmetry:

$$L(\mathbf{k},\mathbf{j}) = L(\mathbf{j},\mathbf{k}).$$

From this it follows, again by inspection, that the integral of the transfer spectrum vanishes, as it must to conserve energy.

Edwards derived this as a self-consistent mean-field solution to the Liouville equation that is associated with the Navier-Stokes equation, and specialised it to the stationary case. Later Orszag [3] derived a similar form by modifying the quasi-normality theory to obtain a closure called the eddy-

damped quasi-normality markovian (or EDQNM) model. Although physically motivated, this was an *ad hoc* procedure and involved an adjustable constant. For this reason it is strictly regarded as a *model* rather than a *theory*. As this closure is much used for practical applications, we write in terms of the energy spectrum  $E(k,t)=4\pi k^2 C(k,t)$  as:

$$T(k,t) = \int_{p+q=k} D(k,p,q)(xy+z^3) E(q,t)[E(p,t)pk^2 - E(k,t)p^3] \frac{dpdq}{pq}, \quad \text{where}$$

$$D(k,p,q) = \frac{1}{\eta(k,t)+\eta(p,t)+\eta(q,t)}, \quad \text{and}$$

$\eta(k,t)$  is the inverse modal response time (equivalent to  $\omega(k,t)$  in the Edwards theory, but determined in a different way). Also  $(xy+z^3)$  is a geometric factor, where  $x$ ,  $y$  and  $z$  are the cosines of the angles of the triangle subtended, respectively, by  $k$ ,  $p$  and  $q$ .

My point here is that Orszag, like many others, followed Kraichnan rather than Edwards and it is clear that you cannot deduce the conservation properties of this formulation by inspection. I should emphasise that the formulation can be shown to be conservative. But it is, in my opinion, much more demanding and complicated than the Edwards form, as I found out when beginning my postgraduate research and I felt obliged to plough my way through it. At one point, Kraichnan acknowledged a personal communication from someone who had drawn his attention to an obscure trigonometrical identity which had proved crucial for his method. Ultimately I found the same identity in one of my old school textbooks [5]. The authors, both masters at Harrow School, had shown some prescience, as they noted that this identity was useful for applications!

During the first part of my research, I had to evaluate integrals which relied on the cancellation of pairs of terms which were separately divergent at the origin in wavenumber. At the time I felt that Kraichnan's way of handling the three

scalar wavenumbers would have been helpful, but I managed it nonetheless in the Edwards formulation. Later on I was to find out, as mentioned in the previous blog, that there were in fact snags to Kraichnan's method too.

In 1990 [4] I wrote about the widespread use of EDQNM in applications. What was true then is probably much more the case today. It seems a pity that someone does not break ranks and employ this useful model closure in the Edwards formulation, rather than make *ad hoc* corrections afterwards for the case of wavenumber triangles with one very small side.

[1] W. David McComb. Homogeneous, Isotropic Turbulence: Phenomenology, Renormalization and Statistical Closures. Oxford University Press, 2014.

[2] S. F. Edwards. The statistical dynamics of homogeneous turbulence. J. Fluid Mech., 18:239, 1964.

[3] S. A. Orszag. Analytical theories of turbulence. J. Fluid Mech., 41:363, 1970.

[4] W. D. McComb. The Physics of Fluid Turbulence. Oxford University Press, 1990.

[5] A. W. Siddons and R. T. Hughes. Trigonometry: Part 2 Algebraic Trigonometry. Cambridge University Press, 1928.