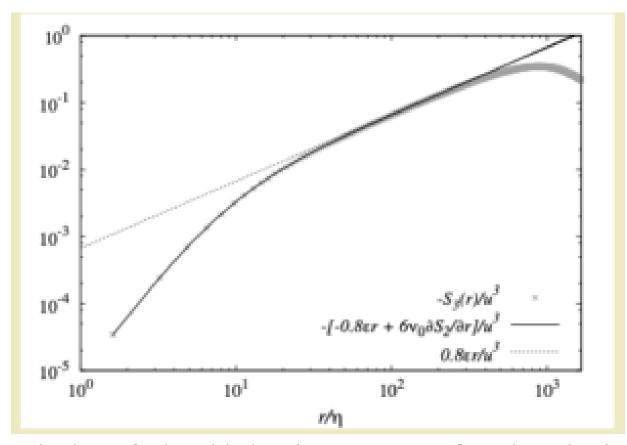
Viscous and forcing corrections to Kolmogorov's '4/5' law.

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The Kolmogorov `4/5′ law for the third-order structure function $S_3(r)$ is widely regarded as the one exact result in turbulence theory. And so it should be: it has a straightforward derivation from the Karman-Howarth equation (KHE), which is an exact energy balance derived from the Navier-Stokes equation. Nevertheless, there is often some confusion around its discussion in the literature. In particular, for stationary isotropic turbulence, there can be confusion about the effects of viscosity (small scales) and forcing (large scales). These aspects have been clarified by McComb et al [1], who used spectral methods to obtain S_2 and S_3 from a direct numerical simulation of the equations of motion.

If we follow the standard treatment (see [2], Section 4.6.2), we may write: \begin{equation} $S_3(r) = -\frac{4}{5} \ r = \frac{1}{100} r + 6 \ r = \frac{100}{100} r + \frac{100}{100} r + 6 \ r = \frac{100}{100} r + \frac$



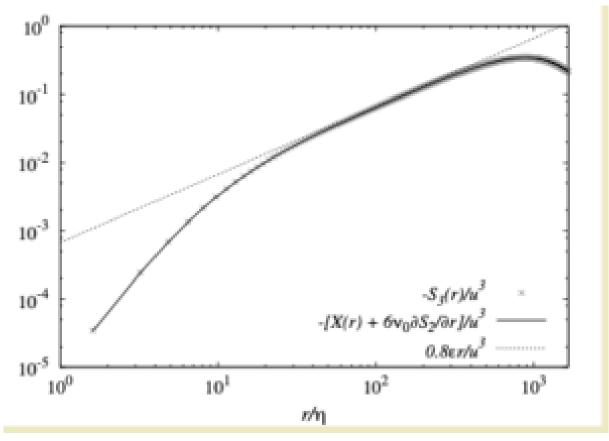
Variation of the third-order structure function showing the effect of viscous corrections.

The results were taken at a Taylor-Reynolds number $R_{\lambda} = 435.2$, and show how the departure from the 4/5 law at the small scales is due to the viscous effects. Clearly there is a range of values of r where the 4/5 law may be regarded as exact, in the ordinary sense appropriate to experimental work. This range of scales is, of course, the inertial range. Note that \det is the Kolmororov length scale.

Presumably the departure from the `4/5' law at the large scales is due to forcing effects, and McComb et al [1] also shed light on this point. They did this by working in spectral space, where stirring forces have been studied since the late 1950s in the context of the statistical theories (e.g Kraichnan, Edwards, Novikov, Herring: see [3] for details) and are correspondingly well understood. They began with the Lin equation: \begin{equation} $frac{\operatorname{E}(k,t)}{\operatorname{Partial t}} = T(k,t) - 2 u k^2E(k,t) + W(k), \end{equation} where in principle the energy and transfer spectra depend on time, whereas the spectrum of the stirring forces $W(k)$ is taken as independent of time in order to ensure ultimate stationarity. Thus we will drop the time dependences hereafter as we will only consider the stationary case.$

We can derive the KHE from this and the result is the usual KHE plus input term \$I(r)\$, defined an by: $\begin{equation}I(r) = \frac{3}{r^3} int 0^r, dy , y^2,$ W(y), $end\{equation\}$ where W(y) is the three-dimensional Fourier transform of the work spectrum \$W(k)\$. By integrating the KHE (as Kolmogorov did in deriving the $^4/5'$ law) we obtain the form for the third-order structure function S 3(r) as: \begin{equation} S 3(r)=X(r) + 6\nu\frac{\partial} S 2}{\partial r},\end{equation}where where \$X(r)\$ is given in terms of the forcing spectrum by: $\sum \{e_i\} X(r) =$ -12r\int $0^{\pm}, dk W(k), \leq x - 3kr \cos$ $kr-(kr)^2 \ kr}{(kr)^5} \ ight].\$

The result of including the effect of forcing is shown in Figure 8 of [1], which is reproduced here below.



Variation of the third-order structure function with scale showing both viscous effects and those due to forcing.

These results are taken from the same simulation as above, and now the contributions from viscous and forcing effects can be seen to account for the departure of S_3 from the 4/5 law at all scales.

[1] W. D. McComb, S. R. Yoffe, M. F. Linkmann, and A. Berera. Spectral analysis of structure functions and their scaling exponents in forced isotropic turbulence. Phys. Rev. E, 90:053010, 2014.

[2] W. David McComb. Homogeneous, Isotropic Turbulence: Phenomenology, Renormalization and Statistical Closures. Oxford University Press, 2014.

[3] W. D. McComb. The Physics of Fluid Turbulence. Oxford University Press, 1990.